## SUPPLEMENTAL MATERIAL

## Derivation of 1D analytical concentration profiles

Steady-state post-enrichment, 1D model
The steady-state equations of change for diffusing species $A, B$, and $C$ are:

$$
\begin{align*}
& \frac{\partial A}{\partial t}=0=D_{A} \frac{\partial^{2} A}{\partial x^{2}}  \tag{1}\\
& \frac{\partial B}{\partial t}=0=D_{B} \frac{\partial^{2} B}{\partial x^{2}}-k_{B} B  \tag{2}\\
& \frac{\partial C}{\partial t}=0=D_{C} \frac{\partial^{2} C}{\partial x^{2}}+k_{B} B \tag{3}
\end{align*}
$$

The boundary conditions are:
$\left.J_{A}\right|_{0}=\left.J_{A}\right|_{L}=J$
$\left.J_{B}\right|_{0}=-\left.J_{A}\right|_{0}=J$
$\left.J_{B}\right|_{L}=0$
$\left.J_{C}\right|_{0}=0$
$\left.J_{C}\right|_{L}=-\left.J_{A}\right|_{0}=J$
where $J$ is the flux. These conditions state that the flux from the reaction surfaces are diffusionlimited (hence they do not contain rate constants and do not depend on reactant concentration) and constant at steady-state. As a basis of calculation, we allow the surface reaction of $B, B(0)$, to reach an arbitrary constant value of $B_{0}$ at steady-state, which will ultimately depend on system parameters such as the reaction rate constants and total amount of protein present in the embryo (see below).

Solving for $A$ :
$\frac{\partial^{2} A}{\partial x^{2}}=0$
$\frac{\partial A}{\partial x}=m$
$A=m x+A_{0}$
where $A_{0}$ is the concentration of $A$ at $x=0$. We see that $A$ is linear, and the flux of $A$ is constant over all $x$ from 0 to $L$.
$J_{A}=-D_{A} \frac{d A}{d x}=-D_{A} m$

Solving for $B$ :

$$
\frac{\partial^{2} B}{\partial x^{2}}=\frac{k_{B}}{D_{B}} B
$$

With steady flux at $x=0$ and zero flux at $x=\mathrm{L}$, we expect the concentration of $B$ to achieve an arbitrary steady value we will denote by the constant $B_{0}$, such that

$$
\begin{equation*}
B=B_{0} \frac{\cosh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}} \tag{10}
\end{equation*}
$$

We see that the flux of $B$ is
$J_{B}=-D_{B} \frac{d B}{d x}=B_{0} \sqrt{k_{B} D_{B}} \frac{\sinh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}$
so the flux of $B$ at 0 is

$$
\left.J_{B}\right|_{0}=B_{0} \sqrt{k_{B} D_{B}} \frac{\sinh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}=B_{0} \sqrt{k_{B} D_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}
$$

and the flux of $B$ at L is

$$
\left.J_{B}\right|_{L}=B_{0} \sqrt{k_{B} D_{B}} \frac{\sinh (0)}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}=0
$$

Solving for $C$ :

$$
\begin{aligned}
& \frac{\partial^{2} C}{\partial x^{2}}=-\frac{k_{B}}{D_{B}} B=-\frac{k_{B}}{D_{B}} B_{0} \frac{\cosh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}} \\
& \frac{\partial C}{\partial x}=\frac{B_{0}}{D_{C}} \sqrt{D_{B} k_{B}} \frac{\sinh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}+c_{1} \\
& C=-B_{0}\left(\frac{D_{B}}{D_{C}}\right) \frac{\left.\cosh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}\left(1-\frac{x}{L}\right.}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}+c_{1} x+c_{2}
\end{aligned}
$$

We can solve for $\mathrm{c}_{1}$ by considering the zero flux condition at $x=0$ :
$\left.J_{C}\right|_{0}=0=-\left.D_{C} \frac{\partial C}{\partial x}\right|_{0}$

$$
\begin{aligned}
& \left.\frac{\partial C}{\partial x}\right|_{0}=0=\frac{B_{0}}{D_{C}} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}+c_{1} \\
& c_{1}=-\frac{B_{0}}{D_{C}} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}
\end{aligned}
$$

Plugging in for $\mathrm{c}_{1}$, we obtain

$$
C=-B_{0}\left(\frac{D_{B}}{D_{C}}\right) \frac{\cosh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}-\frac{B_{0}}{D_{C}} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}} x+c_{2}
$$

The flux of $C$ is

$$
\frac{\partial C}{\partial x}=\frac{B_{0}}{D_{C}} \sqrt{D_{B} k_{B}} \frac{\sinh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}-\frac{B_{0}}{D_{C}} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}
$$

Similarly to the case for $B$, with a steady flux of $C$ at $x=\mathrm{L}$, we expect the concentration of $C$ to attain a steady value at $x=0$ that we will denote $C_{0}$, such that

$$
\begin{aligned}
& C(0)=C_{0}=-B_{0}\left(\frac{D_{B}}{D_{C}}\right)+c_{2} \\
& C_{2}=C_{0}+B_{0}\left(\frac{D_{B}}{D_{C}}\right)
\end{aligned}
$$

Plugging in for $\mathrm{c}_{2}$, we obtain
$C=C_{0}+B_{0}\left(\frac{D_{B}}{D_{C}}\right)\left[1-\frac{\cosh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}\right]-\frac{B_{0}}{D_{C}} \sqrt{k_{B} D_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}} x}$

The flux of $C$ is

$$
J_{C}=-D_{C} \frac{\partial C}{\partial x}=-B_{0} \sqrt{D_{B} k_{B}} \frac{\sinh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}}+B_{0} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}
$$

So the flux at $x=0$ is

$$
\left.J_{C}\right|_{0}=-\left.D_{C} \frac{\partial C}{\partial x}\right|_{0}=-B_{0} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}+B_{0} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}=0
$$

And the flux at $x=\mathrm{L}$ is

$$
\left.J_{C}\right|_{L}=-\left.D_{C} \frac{\partial C}{\partial x}\right|_{L}=B_{0} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}
$$

which, as expected, is equal to the flux of $B$ at $x=0$. Of course, the flux of $B$ at $x=0$ must be equal and opposite of the flux of $A$ at $x=0$. Likewise, the flux of $C$ at $x=\mathrm{L}$ must be equal and opposite to the flux of $A$ at $x=\mathrm{L}$. Hence,

$$
\left.J_{B}\right|_{0}=B_{0} \sqrt{k_{B} D_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}=-\left.J_{A}\right|_{0}=D_{A} m
$$

So

$$
m=\frac{B_{0}}{D_{A}} \sqrt{k_{B} D_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}
$$

We can check the flux at $x=L .$. .

$$
\left.J_{C}\right|_{L}=B_{0} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}=\left.J_{A}\right|_{L}=D_{A} m
$$

And we again find that
$m=\frac{B_{0}}{D_{A}} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}$
Hence,

$$
\begin{equation*}
A=m x+A_{0}=\left(\frac{B_{0}}{D_{A}} \sqrt{D_{B} k_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}\right) x+A_{0} \tag{13}
\end{equation*}
$$

$A_{0}$, which determines the relative concentration of $A$ in the system, is left as a tunable parameter.

The total concentration can then be calculated by

$$
\begin{align*}
& \text { Total }=A+B+C \\
& \text { Total }=A_{0}+C_{0}+B_{0}\left(\frac{D_{B}}{D_{C}}\right) \\
& +\left(\frac{1}{D_{A}}-\frac{1}{D_{C}}\right) B_{0} \sqrt{k_{B} D_{B}} \tanh \sqrt{\frac{k_{B} L^{2}}{D_{B}}} x \\
& +\left(1-\frac{D_{B}}{D_{C}}\right) B_{0} \frac{\cosh \left[\sqrt{\frac{k_{B} L^{2}}{D_{B}}}\left(1-\frac{x}{L}\right)\right]}{\cosh \sqrt{\frac{k_{B} L^{2}}{D_{B}}}} \tag{14}
\end{align*}
$$

In order to test the ability of our model to correctly predict the overall concentration gradient, we measured fluorescence intensity along the $\mathrm{A} / \mathrm{P}$ axis of embryos expressing MEX5::GFP (with background fluorescence subtracted) and used data from the central 40um of the $\mathrm{A} / \mathrm{P}$ axis $(L)$, and used values for the diffusion coefficients $\left(D_{A}, D_{B}, D_{C}\right)$ measured with FCS. Species intensity $\left(A_{0}, B_{0}, C_{0}\right)$ and reaction $\left(k_{B}\right)$ parameters were tuned to minimize squared error. Diffusion coefficients for slow system components ( $B$ and $C$ ) were constants equal to $0.4 \mu^{2} / \mathrm{s}$ and $1.0 \mu \mathrm{~m}^{2} / \mathrm{s}$, respectively, while the diffusion coefficient for the fast species ( $A$ ) was constant and equal to $15 \mu \mathrm{~m}^{2} / \mathrm{s}$ (all of which are within experimental error of our FCS measurements). As
a basis for calculation, the total intensity at the anterior pole ( $x=0 \mu \mathrm{~m}$ ) was normalized to $\operatorname{Total}(0)$ $=5$ (a.u.).

The rate constants used in our modeling are modified such that $k_{\text {cytoplasmic }}=k_{\text {phys }} k_{I}[P]$ and $k_{\text {surface }}=k_{\text {phys }} k_{I} \frac{E_{\text {tot }}}{K_{M}}$ (assuming first order surface reactions), where $k_{\text {phys }}$ is the physiological rate constant, $k_{I}=\frac{[S]}{S}$ is an (unknown) constant relating fluorescence intensity (S) to cytoplasmic concentration ([S]), $[P]$ is the (unknown) concentration of the putative cytoplasmic phosphatase, $E_{\text {tot }}$ is the (unknown) total surface concentration of enzyme on the reaction surface, $K_{M}$ is the Michaelis constant, and $k_{\text {phys }}$ is the physiological rate constant. The modified rate constants estimated from the modeling were $k_{A}=0.065 \mu \mathrm{~m} / \mathrm{s}, k_{B}=0.0025 \mathrm{~s}^{-1}$, and $k_{C}=0.2 \mu \mathrm{~m} / \mathrm{s}$, and the concentration parameters (a.u.) were estimated to be $A_{0}=1, B_{0}=2.1, C_{0}=1.9$.

## FEMLAB

Multi-dimensional solutions to the partial differential equations were found using the finite element method (FEM) via COMSOL 3.5a (FEMLAB) software. FEM analysis is a useful tool for calculating solutions to complicated systems of equations over complex geometries including intricate biological problems (Reddy, 1993; Sun et al., 2009). Briefly, the finite element method calculates the solution to complicated systems by breaking the geometry of a given system into subdomains, or finite elements. Solutions to each finite element are calculated by approximating solutions to the partial differential equations over that element as a linear sum of algebraic polynomials where the undetermined coefficients of these polynomials are valued according to the governing partial differential equations. The finite elements themselves are not fixed throughout solving the governing system of equations but may change to reduce the error
in approximation over a given element and maintain continuity of the solution over all elements (Reddy, 1993).

Here, we used a simple oval geometry (major and minor axes of $50 \mu \mathrm{~m}$ and $30 \mu \mathrm{~m}$, respectively) centered about the origin, consisting of two spatially identical and opposite reactive surfaces, the anterior and posterior, to represent the dividing embryo. The $A$ to $B$ reactions were allowed to occur only on the anterior surface, the $B$ to $C$ reaction again occurred everywhere within the cytoplasmic region and the $C$ to $A$ reaction was allowed to occur only on the posterior surface. All reactions rates used to calculate the multi-dimensional solutions were first order. Diffusion coefficients for slow system components ( $B$ and $C$ ) were constants equal to $0.4 \mu \mathrm{~m}^{2} / \mathrm{s}$ and $1.0 \mu \mathrm{~m}^{2} / \mathrm{s}$, respectively, while the diffusion coefficient for the fast species ( $A$ ) was constant and equal to $15 \mu \mathrm{~m}^{2} / \mathrm{s}$ (all of which are within experimental error of our FCS measurements). As a basis for calculation, the concentration of $A$ at the anterior pole (far left anterior node) was considered a known quantity and set equal to 1 (a.u.). The modified rate constants (see above) used in the modeling were $k_{A}=0.02 \mu \mathrm{~m} / \mathrm{s}, k_{B}=0.001 \mathrm{~s}^{-1}$ and $k_{C}=0.5 \mu \mathrm{~m} / \mathrm{s}$.

## Analytical concentration profiles with reactive species B

The governing equations are identical to those for the non-reacting $B$ species, however the zeroflux boundary condition at $x=L$ must be allowed to take some finite value, $\beta$. Thus, the general equation for species $B$ becomes
$B=a \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+b \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)$

We specify that $B$ attains for steady-state value at $x=0$ such that
$B(0)=B_{0}=a+b$
$b=B_{0}-a$
$B=a \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\left(B_{0}-a\right) \exp \left(-\sqrt{\frac{k_{B}}{D_{B}} x}\right)$
$B=a \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-a \exp \left(-\sqrt{\frac{k_{B}}{D_{B}} x} x\right)+B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}} x}\right)$
$B=a\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]+B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)$
and
$\frac{d B}{d x}=\sqrt{\frac{k_{B}}{D_{B}}} a \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\sqrt{\frac{k_{B}}{D_{B}}}\left(B_{0}-a\right) \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)$
$\frac{d B}{d x}=\sqrt{\frac{k_{B}}{D_{B}}} a \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\sqrt{\frac{k_{B}}{D_{B}}} a \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\sqrt{\frac{k_{B}}{D_{B}}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)$
$\frac{d B}{d x}=a \sqrt{\frac{k_{B}}{D_{B}}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]-\sqrt{\frac{k_{B}}{D_{B}}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)$
We allow the flux of $B$ to attain some finite value, $\beta$, at steady-state such that

$$
\begin{aligned}
& J_{B}=-\left.D_{B} \frac{d B}{d x}\right|_{x=L}=\beta=-\sqrt{k_{B} D_{B}} a \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\sqrt{k_{B} D_{B}}\left(B_{0}-a\right) \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right) \\
& \beta=B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-a \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-a \sqrt{k_{B} D_{B}} \exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right) \\
& \beta=B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-a \sqrt{k_{B} D_{B}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right] \\
& a=\frac{B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta}{\sqrt{k_{B} D_{B}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \\
& B=\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{\sqrt{k_{B} D_{B}}}+\exp ^{\left.\left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}+B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right) \\
& \frac{d B}{d x}=\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{D_{B}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]-\sqrt{\frac{k_{B}}{D_{B}}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)
\end{aligned}
$$

where $\beta$ is a constant. Thus, the flux of B is given by

$$
J_{B}=-D_{B} \frac{d B}{d x}=\left[\beta-B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right] \frac{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}}\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}+\sqrt{k_{B} D_{B}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)
$$

We can use the explicit expression for $B$ to solve for $C$ by integration

$$
\left.\left.\begin{array}{l}
\frac{d^{2} C}{d x^{2}}=-\frac{k_{B}}{D_{C}} B=-\frac{k_{B}}{D_{C}} \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{\sqrt{k_{B} D_{B}}} \underset{\left.\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}{ }-\frac{k_{B}}{D_{C}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right) \\
\frac{d C}{d x}=-\frac{\left[B_{0} \sqrt{k_{B_{B} D_{B}}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{D_{C}}+\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)+c_{1} \\
\exp
\end{array}\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right] \quad .
$$

The flux of $C$ at $x=0$ is specified to equal 0 such that

$$
\begin{aligned}
& \left.\frac{d C}{d x}\right|_{x=0}=0=-\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \frac{2}{D_{C}}+\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0}+c_{1} \\
& c_{1} \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \frac{2}{D_{C}}-\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0} \\
& \frac{d C}{d x}=-\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{D_{C}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right] \frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}} x}\right) \\
& +\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \frac{2}{D_{C}}-\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0}
\end{aligned}
$$

$$
\begin{aligned}
& C=c_{2}-\sqrt{\frac{D_{B}}{k_{B}}} \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{D_{C}} \frac{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \\
& -\frac{D_{B}}{D_{C}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\left\{\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \frac{2}{D_{C}}-\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0}\right\} x
\end{aligned}
$$

Finally, we specify that $C$ attains a constant value, $C_{0}$, at $x=0$ :
$C(0)=C_{0}=c_{2}-\frac{D_{B}}{D_{C}} B_{0}$
$c_{2}=C_{0}+\frac{D_{B}}{D_{C}} B_{0}$
such that

$$
\begin{aligned}
& C=C_{0}+\frac{D_{B}}{D_{C}} B_{0}-\sqrt{\frac{D_{B}}{k_{B}}} \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{D_{C}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right] \\
& -\frac{D_{B}}{D_{C}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\left\{\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \frac{2}{D_{C}}-\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0}\right\} x
\end{aligned}
$$

and thus

$$
\begin{aligned}
& \frac{d C}{d x}=-\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{D_{C}}\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right] \\
& +\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left.\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]\right]} \frac{2}{D_{C}}-\frac{\sqrt{k_{B} D_{B}}}{D_{C}} B_{0}
\end{aligned}
$$

Therefore, the flux of $C$ is given by

$$
\begin{aligned}
& J_{C}=-D_{C} \frac{d C}{d x}=\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right] \frac{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \\
& -\sqrt{k_{B} D_{B}} B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)-2 \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]+\sqrt{k_{B} D_{B}} B_{0}}
\end{aligned}
$$

As expected, the flux of $C$ at $x=0$ is

$$
\begin{aligned}
& \left.J_{C}\right|_{x=0}=-\left.D_{C} \frac{d C}{d x}\right|_{x=0}=\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right] \frac{[1+1]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]} \\
& -\sqrt{k_{B} D_{B}} B_{0}-2 \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}+\sqrt{k_{B} D_{B}} B_{0}=0
\end{aligned}
$$

Similarly, the flux of $C$ at $x=\mathrm{L}$ is

$$
\left.J_{C}\right|_{x=L}=-\left.D_{C} \frac{d C}{d x}\right|_{x=L}=\sqrt{k_{B} D_{B}} B_{0}-2 \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}-\beta
$$

We can verify that the flux of $B$ at $x=0$ equals the sum of the flux of $C$ at $x=L$ and the flux of $B$ at $x=L:$

$$
J_{B}(0)=\sqrt{k_{B} D_{B}} B_{0}-2 \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}
$$

$$
J_{B}(L)=\beta
$$

Finally, the flux of $A$ at $x=0$ must be equal and opposite to the flux of $B$ at $x=0$, such that
$=\sqrt{k_{B} D_{B}} B_{0}-2 \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}=D_{A} m$
$m=\frac{\sqrt{k_{B} D_{B}}}{D_{A}} B_{0}-\frac{2}{D_{A}} \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}$

And the concentration profile of $A$ is given by

$$
A=A_{0}+\left\{\frac{\sqrt{k_{B} D_{B}}}{D_{A}} B_{0}-\frac{2}{D_{A}}\left[\operatorname{Bexp}\left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]\right\}
$$

The (constant) flux of $A$ is given by

$$
J_{A}=-D_{A} \frac{d A}{d x}=\frac{\sqrt{k_{B} D_{B}}}{D_{A}} B_{0}-\frac{2}{D_{A}} \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}
$$

The total concentration profile is then given by
Total $=A+B+C$

$$
\begin{align*}
& =A_{0}+C_{0}+\frac{D_{B}}{D_{C}} B_{0}+\left(1-\frac{D_{B}}{D_{C}}\right) B_{0} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right) \\
& \left.+\left(1-\frac{D_{B}}{D_{C}}\right) \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]\left[\exp \left(\sqrt{\frac{k_{B}}{D_{B}}} x\right)-\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} x\right)\right]}{\sqrt{k_{B} D_{B}}}\left(\sqrt{\frac{k_{B}}{D_{B}}} L\right)+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]  \tag{15}\\
& +\left(\frac{1}{D_{C}}-\frac{1}{D_{A}}\right)\left\{2 \frac{\left[B_{0} \sqrt{k_{B} D_{B}} \exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)-\beta\right]}{\left[\exp \sqrt{\frac{k_{B}}{D_{B}}} L+\exp \left(-\sqrt{\frac{k_{B}}{D_{B}}} L\right)\right]}-\sqrt{k_{B} D_{B}} B_{0}\right\} x
\end{align*}
$$

